

Lecture 9: Wave - mean flow interactions.

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Course GFD M2 MOCIS

Interaction Rossby wave - mean current

Mean fields, fluctuation fields, and their interactions

Wave pseudomomentum

Rossby waves over a constant mean flow

Wave-mean interaction

Critical levels

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Gravity-wave meanflow interactions

Equations of motion

$$\begin{aligned}u_t + uu_x + vv_y - f(y)v + \frac{1}{\rho}p_x &= 0, \\v_t + uv_x + vv_y + f(y)u + \frac{1}{\rho}p_y &= 0, \\u_x + v_y &= 0.\end{aligned}\tag{1}$$

Typical configuration- zonal channel: $-D \leq y \leq D$,
b.c.: $v|_{\pm D} = 0$.

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Decomposition: zonal mean/fluctuations

Zonal mean and fluctuations:

Zonal mean:

$$\bar{A} = \lim_{L \rightarrow \infty} \frac{1}{2L} \int_{-L}^{+L} dx A(x, y, t)$$

$$\bar{A} = \frac{1}{2L} \int_{-L}^{+L} dx A(x, y, t), \text{ si } A(x + 2L) = A(x). \quad (2)$$

Properties:

$$A = \bar{A} + \tilde{A}; \quad \tilde{A} \equiv A - \bar{A}; \quad \overline{A_x} \equiv 0; \quad \overline{\tilde{A}} \equiv 0.$$

$$\overline{\tilde{A}} = \bar{A}; \quad \overline{\tilde{A}\tilde{B}} \equiv 0.$$

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Mean barotropic zonal current in a channel

$$(u, v) = (\bar{u} + \tilde{u}, \bar{v} + \tilde{v})$$
$$\overline{u_x + v_y} = 0, \Rightarrow \bar{v}_y = \bar{v}_y = 0, \Rightarrow \bar{v} = 0. \quad (3)$$
$$\tilde{u}_x + \tilde{v}_y = 0.$$

Therefore:

$$\overline{u_t + uu_x + vu_y - f(y)v + \frac{1}{\rho}p_x} = 0 \Rightarrow$$
$$\bar{u}_t + \overline{\tilde{v}\tilde{u}_y} = \bar{u}_t + \overline{(\tilde{v}\tilde{u})_y} = 0. \quad (4)$$

Physical meaning: **change of the mean zonal momentum**
 \leftrightarrow **momentum flux of the fluctuations.**

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Zonal flow + waves of small amplitude

Linearisation:

$\bar{u} = U(y)$, fluctuations: $\|\tilde{u}, \tilde{v}\| \ll 1$.

Absolute (potential) vorticity:

$$q = v_x - u_y + f(y) \equiv Q(y) + \tilde{q}. \quad (5)$$

$$Q(y) = f_0 + \beta y - U'(y), \quad \tilde{q} = \tilde{v}_x - \tilde{u}_y. \quad (6)$$

Dynamics of the mean flow:

$$\bar{u}_t = -\overline{(\tilde{v}\tilde{u})}_y = \overline{(\tilde{q}\tilde{v})} \quad (7)$$

Change of zonal flow \leftrightarrow potential vorticity flux

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Pseudomomentum

Linearised vorticity equation :

$$\frac{D\tilde{q}}{Dt} + \tilde{v}Q'(y) = 0, \quad \frac{D}{Dt} = \partial_t + U(y)\partial_x. \quad (8)$$

Observation:

$$\overline{-\frac{\tilde{q}}{Q'(y)} [\tilde{q}_t + U(y)\tilde{q}_x + \tilde{v}Q'(y)]} = 0, \Rightarrow \quad (9)$$

Pseudomomentum:

$$\left[-\frac{\tilde{q}^2}{2Q'(y)} \right]_t - \tilde{q}\tilde{v} = 0. \quad (10)$$

pseudomomentum: $p(y, t) \equiv -\frac{\tilde{q}^2}{2Q'(y)} \Rightarrow$ zonal flow:

$$\bar{u}_t = p_t + \dots \quad (11)$$

Alternative formulation:

$$\tilde{v} = \frac{D\tilde{Y}}{Dt}, \Rightarrow \quad (12)$$

$$\frac{D\tilde{q}}{Dt} + \frac{D\tilde{Y}}{Dt} Q'(y) = 0, \Leftrightarrow \frac{D}{Dt} [\tilde{q} + \tilde{Y} Q'(y)] = 0, \Rightarrow \quad (13)$$

$$\tilde{Y} = -\frac{\tilde{q}}{Q'(y)} + \text{const}, \Rightarrow \quad (14)$$

$$p = -Q'(y) \frac{\overline{\tilde{Y}^2}}{2} = \frac{\overline{\tilde{Y}\tilde{q}}}{2} \quad (15)$$

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Rossby waves over a mean current

Linear waves, dispersion:

Wave solutions $\propto e^{i(kx+ly-\omega t)}$. **Absolute** frequency:

$$\omega = Uk - \beta \frac{k}{k^2 + l^2}. \quad (16)$$

Intrinsic frequency: $\hat{\omega} = -\beta \frac{k}{k^2 + l^2}$. Phase velocity:

$\vec{c} = \frac{\omega}{k^2 + l^2} \vec{k}$. Group velocity:

$$\vec{c}_g = (u_g, v_g) = (\partial_k \omega, \partial_l \omega) = \left(U + \beta \frac{k^2 - l^2}{k^4}, 2\beta \frac{kl}{k^4} \right). \quad (17)$$

Momentum flux produced by waves:

$$\overline{\tilde{u}\tilde{v}} = -\frac{l}{k} \overline{\tilde{v}^2} \equiv -\frac{kl}{k^2}, \Rightarrow$$

$$-\text{sign}(\overline{\tilde{u}\tilde{v}}) = \text{sign}kl = \text{sign}v_g.$$

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Pseudomomentum:

$$U(y) = \text{const} \Rightarrow \beta > 0, \Rightarrow p = -\beta \frac{\overline{\tilde{Y}^2}}{2} \leq 0.$$

Physical meaning of pseudomomentum for a wave system:

Mean energy of a wave system:

$$\overline{E} = \frac{1}{2} \left(\overline{\tilde{u}^2 + \tilde{v}^2} \right) = \frac{\hat{\omega}}{k} p, \leftrightarrow p = \frac{k}{\hat{\omega}} \overline{E}, \quad (18)$$

because $\tilde{u} = -\frac{1}{k} \tilde{v}$ and $-\hat{\omega}^2 \overline{\tilde{Y}^2} = \overline{\tilde{v}^2}$.

Conservation of pseudomomentum:

$$\overline{\tilde{u}\tilde{v}} = -\frac{1}{k} \overline{\tilde{v}^2} = +\frac{l\hat{\omega}^2}{k} \overline{\tilde{Y}^2} = p v_g, \Rightarrow \quad (19)$$

$$p_t + (p v_g)_y = 0. \quad (20)$$

Generation of Rossby waves by undulating boundary

Ondulating boundary

Channel with a constant mean flow and undulating boundary:

$$-D + h_0 \cos k_0 x \leq y \leq +D, \quad |h_0| \ll 1, \quad k_0 > 0. \Rightarrow (21)$$

Linearised b.c.:

$$\tilde{v}|_{y=-D} = \frac{D(h_0 \cos kx)}{Dt} = U(h_0 \cos k_0 x)_x \quad (22)$$

$$\Leftrightarrow \tilde{Y}|_{y=-D} = h_0 \cos k_0 x.$$

Emitted waves:

$k = k_0, \omega = Uk_0 + \hat{\omega}$ - fixed by the b. c. \Rightarrow

$$\hat{\omega} = -Uk_0, \Rightarrow U = \frac{\beta}{k_0^2 + l^2}, \Rightarrow l^2 = \frac{\beta}{U} - k_0^2. \quad (23)$$

- **Propagative** waves (Charney - Drazin):

$$l^2 > 0, \Rightarrow 0 < U < \frac{\beta}{k_0^2} \Rightarrow \quad (24)$$

- **Evanescent** (trapped) waves:

$$l^2 < 0.$$

Propagation:

Out of the source (to the North):

$$v_g > 0, \Rightarrow l > 0 (k_0 > 0) \Rightarrow l = \sqrt{\frac{\beta}{U} - k_0^2}$$

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Charney - Drazin theorem

Pseudomomentum:

$$\tilde{Y} = h_0 \cos(k_0 x + ly), \Rightarrow p = -\frac{\beta}{4} h_0^2 = \text{const} \quad (25)$$

Stationary vs non-stationary waves:

- ▶ **Stationary** configuration:

$$p = \text{const} \Rightarrow p_t = 0 \Rightarrow \bar{u}_t = 0.$$

stationary waves can not change the mean flow
(Charney-Drazin)

- ▶ **Non-stationary** configuration : Example: h_0 increases from zero to $h_0 \Rightarrow$ **net result:**

$$\bar{u} = U + p = U - \frac{\beta}{4} h_0^2 \rightarrow$$

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Critical levels

Mean flow with a shear:

$U = U(y)$, **critical level**: if $\exists y_c : U(y_c) = 0$.

Stationary perturbations: (for ex. generated by
ondulating boundary)

Streamfunction of perturbations $\tilde{\psi} = \hat{\psi}(y)e^{ikx}$, \Rightarrow potential
vorticity:

$$\tilde{q} = \left(\hat{\psi}''(y) - k^2 \hat{\psi}(y) \right) e^{ikx}.$$

Vorticity equation:

$$U(y)\hat{\psi}''(y) + \left[\beta - U''(y) - k^2 U(y) \right] \hat{\psi}(y) = 0 \quad (26)$$

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Ray tracing and critical levels

WKB method:

Solution is represented as:

$$\tilde{\psi} = \hat{\psi}(y) e^{ikx} e^{i \int_{-D}^y dy' l(y')}, \quad (27)$$

with

$$l(y) = \sqrt{\frac{\beta}{U} - k^2} \quad (28)$$

- **slow (adiabatic) change of wave parameters.**

Condition of applicability of the method: $\left(\frac{U'(y)}{U(y)}\right)^2 \ll l^2$.

Critical level:

$$y \rightarrow y_c, \Rightarrow U(y) \rightarrow 0, l \rightarrow \infty.$$

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In the vicinity of the critical level:

$$U(y) \propto y_c - y, \Rightarrow I(y) \propto (y_c - y)^{-\frac{1}{2}}.$$

Hence:

$$v_g = 2\beta \frac{kl}{(k^2 + l^2)^2} \approx l^{-3} \propto (y_c - y)^{\frac{3}{2}} \rightarrow 0$$

wave-front stops.

Estimate of **approach time**:

$$t = \int dt' = \int \frac{dy'}{v_g(y')} \propto \int \frac{dy'}{(y_c - y')^{\frac{3}{2}}} \propto (y_c - y)^{-\frac{1}{2}} \rightarrow \infty \quad (29)$$

- accumulation below $y_c \Rightarrow$ **critical layer.**

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Breaking near the critical level:

Potential vorticity:

$$q = f_0 + \beta y + \beta \tilde{Y} + \dots \Rightarrow$$

$$q_y = \beta (1 + \tilde{Y}_y)$$

- ▶ Quasi-linear waves: $|\tilde{Y}_y| \ll 1$,
- ▶ Breaking: $|\tilde{Y}_y| \gg 1$.

Estimate: $\overline{\tilde{Y}_y^2} \approx l^2 \overline{\tilde{Y}^2} \approx l^2 p$. Conservation of pseudomomentum for stationary waves:

$$(p v_g)_y = 0 \Rightarrow p(y) \propto \frac{1}{v_g} \Rightarrow \overline{\tilde{Y}_y^2} \propto \frac{l^2}{v_g} \propto l^5 \propto (y_c - y)^{-\frac{5}{2}} \Rightarrow$$

Breaking approaching y_c

Exercise

System: stratified fluid with Brunt - Väisälä frequency $N = \text{const}$, in the vertical plane $x - z$ without rotation in Boussinesq approximation.

- ▶ Consider internal gravity waves over a mean flow $\vec{U} = U\hat{x}$, $U = \text{const}$. Determine absolute and intrinsic frequencies and group velocity,
- ▶ Introduce the **pseudomomentum** $p = -\overline{\tilde{Z}(\tilde{u}_z - \tilde{w}_x)}$, where u, w - horizontal and vertical velocity components, \tilde{Z} - vertical displacements of fluid parcels, $\overline{\dots}$ - **horizontal mean**, and demonstrate that $p_t + \overline{\tilde{u}\tilde{w}_z} = 0$,
- ▶ Demonstrate that $p = k\frac{E}{\omega}$, where k is the horizontal wavenumber, $E = \frac{1}{2}\overline{\tilde{u}^2 + \tilde{w}^2} + \frac{\tilde{\sigma}^2}{N^2}$ - wave energy. Demonstrate that the pseudomomentum flux is $p w_g = \overline{\tilde{u}\tilde{w}}$,

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- ▶ Demonstrate that $\bar{u}_t = \overline{(\tilde{u}\tilde{w})}_z$, and that for the waves of small amplitude $\bar{u}_t = \rho_t + \dots$,
- ▶ Consider **mountain waves** produced by the relief $h(x) = H_0 \cos k_0 x$. Demonstrate that the propagative waves (in vertical direction) exist if $0 < U^2 < \frac{N^2}{k_0^2}$. Determine corresponding w_g and ρ ,
- ▶ Calculate the **mountain drag** $D = -\overline{(\tilde{u}\tilde{w})}$.